

# $W +$ jets and heavy flavour production

All-D0 meeting, July 23rd 2004

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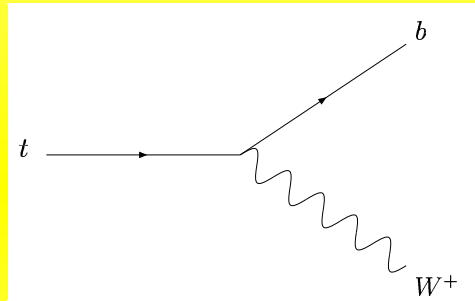
# Outline

- The MCFM Monte Carlo
  - ★ overview of the program
  - ★ implementation of vector boson + jets processes
- NLO predictions for  $W + 2 \text{ jets}$  and  $Wb\bar{b}$ 
  - ★ basic cross-sections
  - ★ theoretical similarities between the two processes
  - ★ effect of NLO on distributions
- Single-tagged events
- Summary

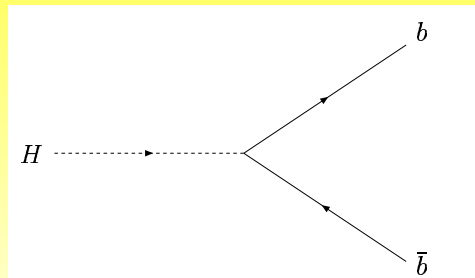
# Heavy flavour as a background

- Events containing jets that are heavy-quark tagged are important for understanding both old and new physics:

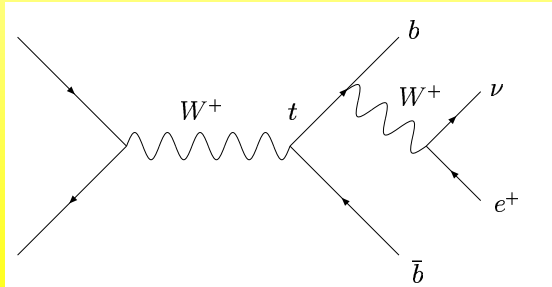
- ★ Top decays  $t \rightarrow W + b$



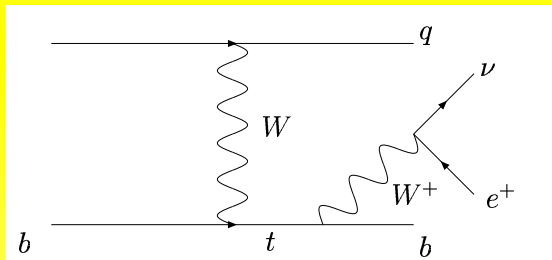
- ★ Much new physics couples preferentially to massive quarks, for instance a light Higgs with  $m_H < 140$  GeV decaying to  $b\bar{b}$



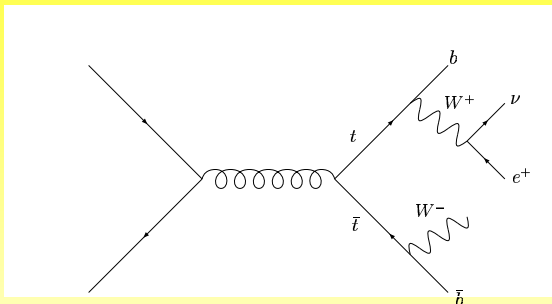
# Top processes



→ 2 jets, both  $b$ 's



→ 2 jets, only one is a  $b$



→  $\geq 2$  jets, two are  $b$ 's

# Overview of MCFM

# MCFM overview

JC and R.K. Ellis

- Parton level cross-sections predicted to NLO in  $\alpha_S$

$p\bar{p} \rightarrow W^\pm / Z$	$p\bar{p} \rightarrow W^+ + W^-$
$p\bar{p} \rightarrow W^\pm + Z$	$p\bar{p} \rightarrow Z + Z$
$p\bar{p} \rightarrow W^\pm + \gamma$	$p\bar{p} \rightarrow W^\pm / Z + H$
$p\bar{p} \rightarrow W^\pm + g^* (\rightarrow b\bar{b})$	$p\bar{p} \rightarrow Z b\bar{b}$
$p\bar{p} \rightarrow W^\pm / Z + 1 \text{ jet}$	$p\bar{p} \rightarrow W^\pm / Z + 2 \text{ jets}$
$p\bar{p}(gg) \rightarrow H$	$p\bar{p}(gg) \rightarrow H + 1 \text{ jet}$
$p\bar{p}(VV) \rightarrow H + 2 \text{ jets}$	

- ⊖ low particle multiplicity (no showering)
- ⊖ no hadronization
- ⊖ hard to model detector effects
- ⊕ less sensitivity to  $\mu_R, \mu_F$
- ⊕ rates are better normalized
- ⊕ fully differential distributions

# *MCFM Information*

- Version 3.4.5 available at:

<http://mcfm.fnal.gov>

- Improvements over previous releases:

- ★ more processes
- ★ better user interface
- ★ support for PDFLIB, Les Houches PDF accord  
( $\longrightarrow$  PDF uncertainties)
- ★ ntuples as well as histograms
- ★ unweighted events
- ★ Pythia/Les Houches generator interface (LO)
- ★ 'Behind-the-scenes' efficiency

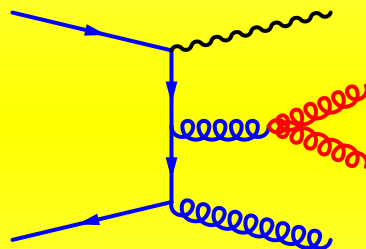
- Coming attractions:

- ★ more processes ( $Z + b$ , single top, ...)
- ★ separate variation of factorization and renormalization scales

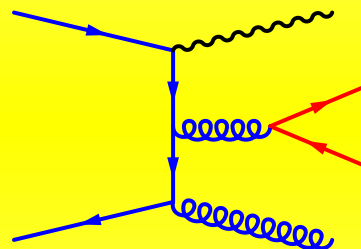
# Vector boson + jets in MCFM

- Many diagrams, sensitive to all parton PDF's
- NLO corrections are separated into two classes:
- REAL extra parton radiation, e.g. ( $W/Z + 2$  jets)

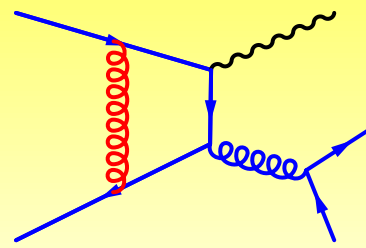
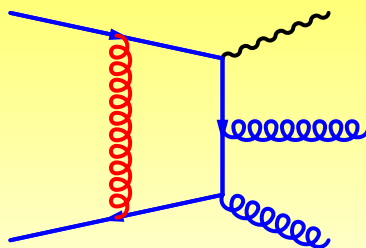
soft gluon



collinear quarks



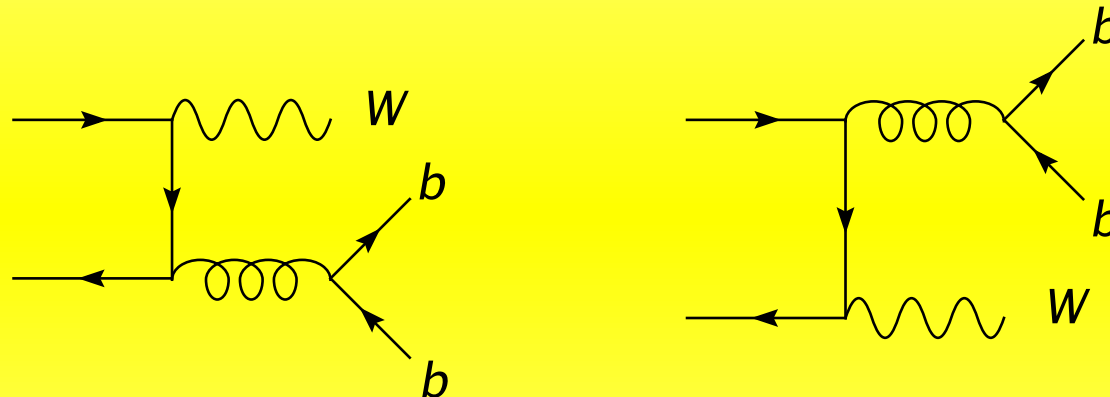
- VIRTUAL loop diagrams:





# Vector boson + heavy flavour in MCFM

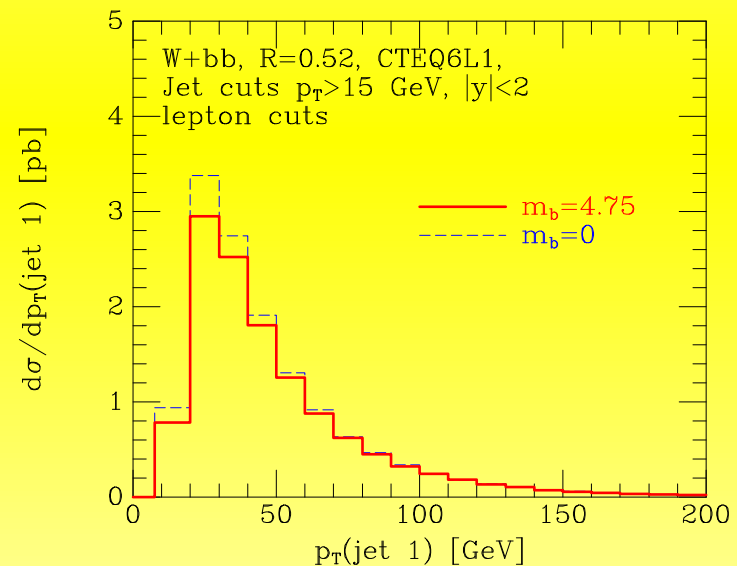
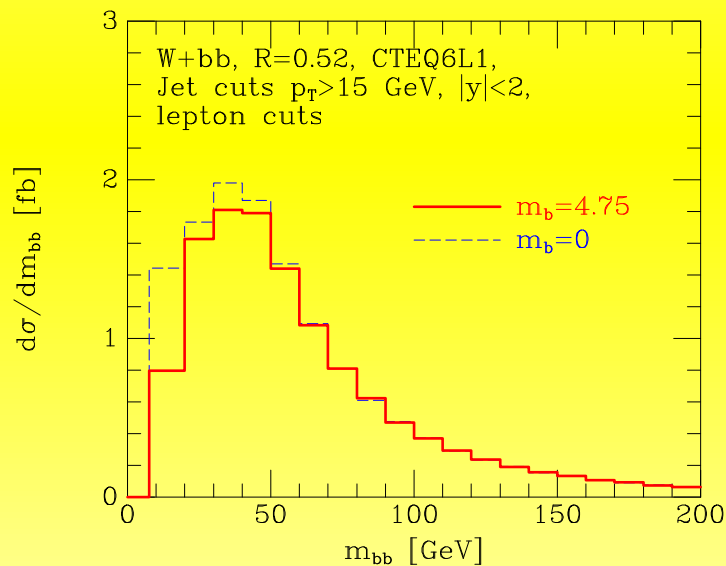
- In lowest order  $b$ -quark pairs are produced in association with  $W$ 's by gluon splitting alone:



- Beyond LO, the  $b$ -quark is treated as a massless particle in MCFM
  - ★ a finite cross-section requires a cut on the  $b$ -quark  $p_T$
  - ★ this means that this calculation is not suitable for estimating the rate with only a single  $b$  tag

# Mass effects

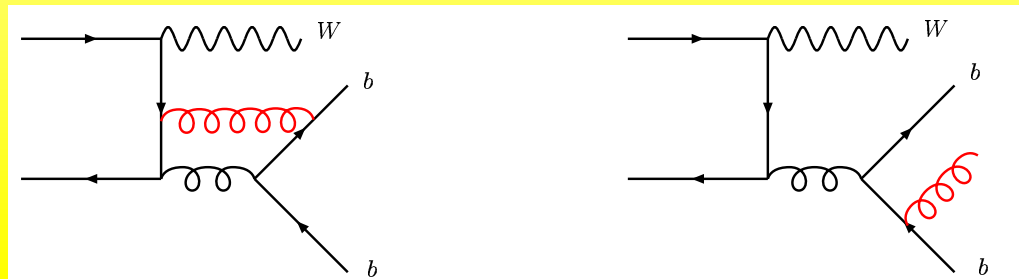
- Lowest order comparison of dijet mass and leading jet  $p_T$  distributions for  $Wb\bar{b}$ 
  - ★  $m_b = 4.75$  GeV (lowest order only)
  - ★  $m_b = 0$  (can be calculated to NLO)



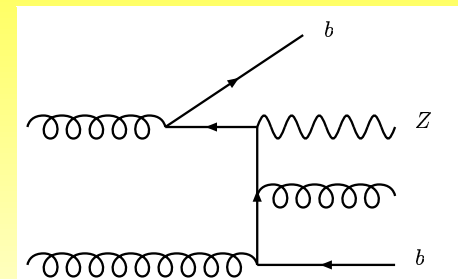
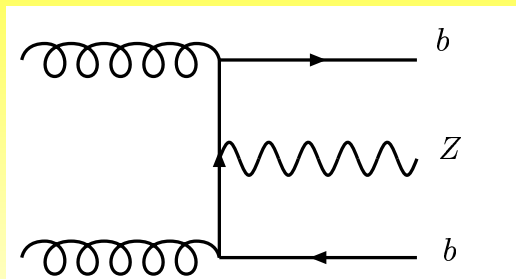
- Overall the cross section decreases by approximately 10% when including the mass. Kinematic distributions are not much affected away from regions of low  $p_T(b)$ .

# Heavy flavour beyond lowest order

- At NLO, the simple kinematics can be altered:

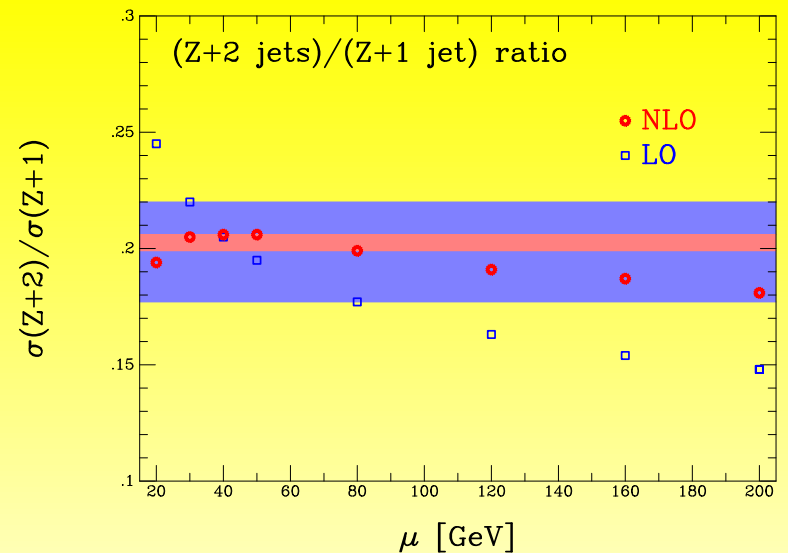
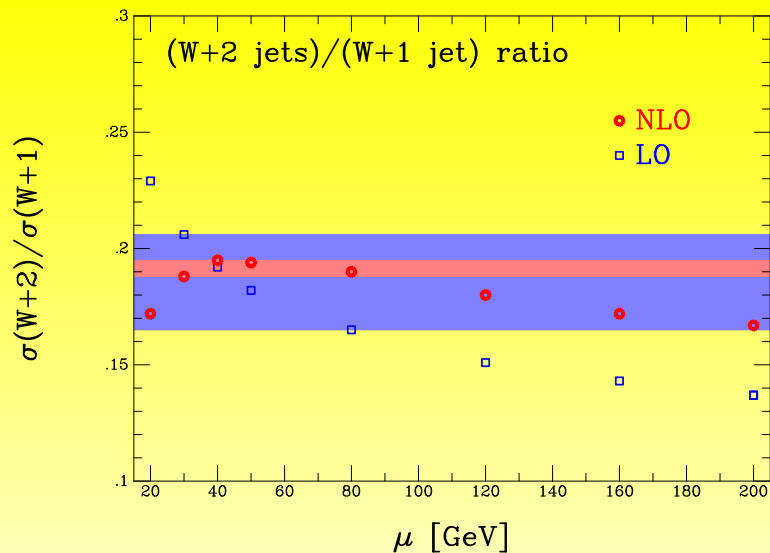


- For heavy flavour production in association with a  $Z$ , the  $b$ -quarks do not have to be produced by gluon splitting. Beyond LO, the difference is further magnified.



# $W/Z + \text{jet}$ cross-sections

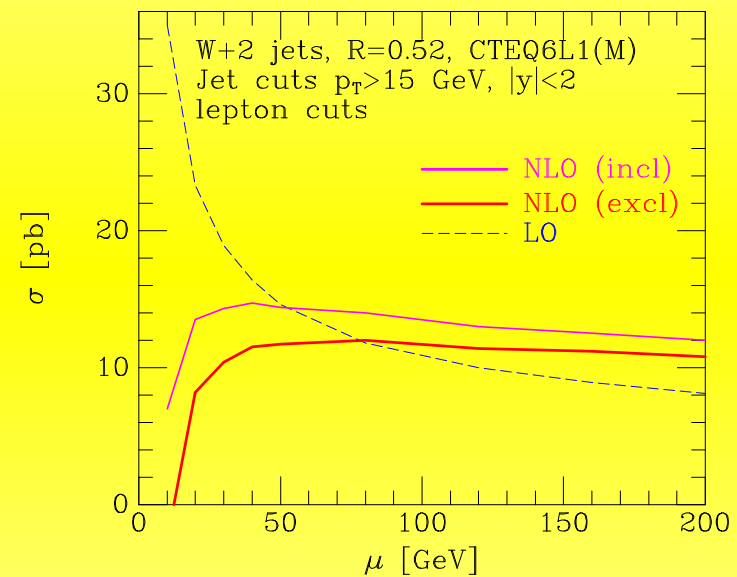
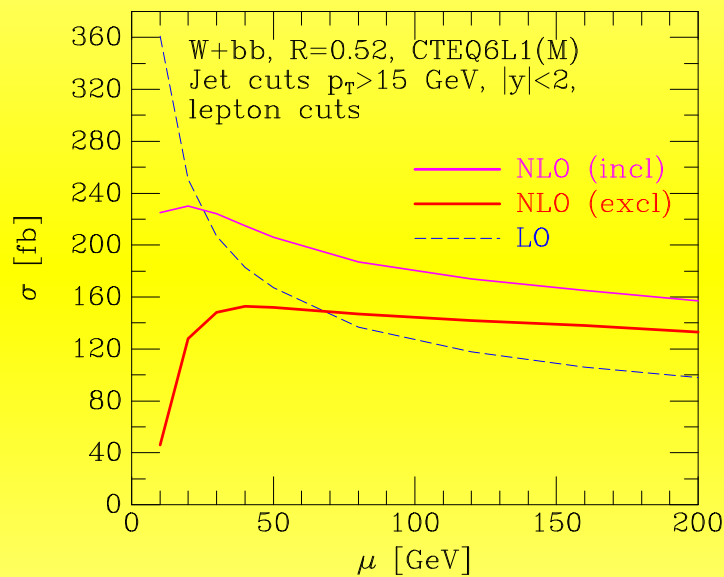
- The  $W/Z + 2$  jet cross-section has only recently been calculated at NLO and should provide an interesting test of QCD (cf. many Run I studies using the  $W/Z + 1$  jet calculation in DYRAD)
- For instance, the theoretical prediction for the number of events containing 2 jets divided by the number containing only 1 is greatly improved.



# Effect of NLO corrections on $Wb\bar{b}$ and $W + 2$ jet rates

# Scale dependence

- Usual scale dependence, much reduced at NLO. Corrections are modest at typical scales,  $\mu \sim M_W$ .

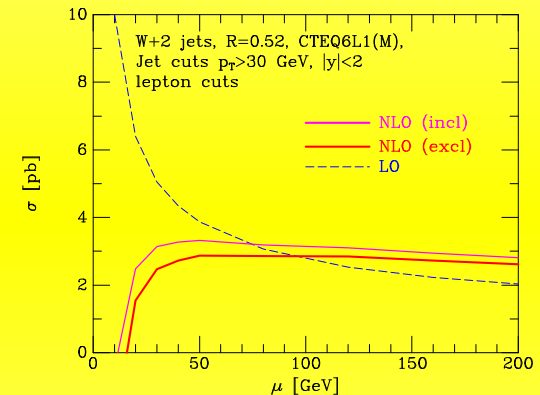
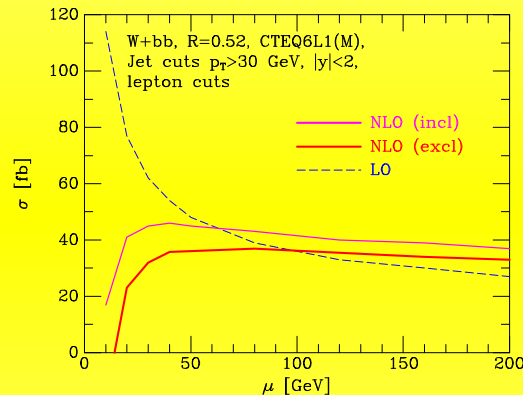


- **Exclusive** cross-sections stable over a large range of scales.
- **Inclusive** result (allows  $Wb\bar{b}j$ ,  $W + 3$  jet configurations) shows more scale dependence, as expected (but still better than LO).

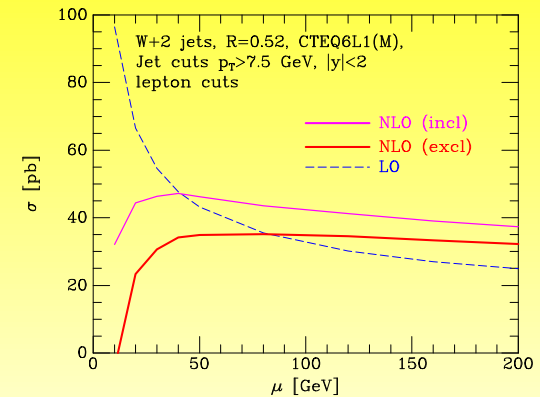
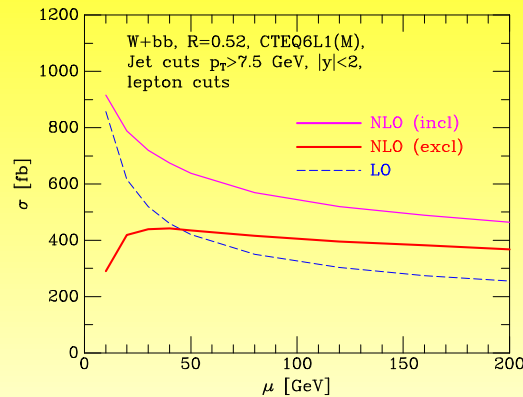
# Jet $p_T$ dependence

- Increasing the minimum jet  $p_T$  reduces the 3 jet contribution compared to the 2 jet one, so the behaviour of the inclusive cross-section improves.

$p_T(\text{jet}) > 30 \text{ GeV}$



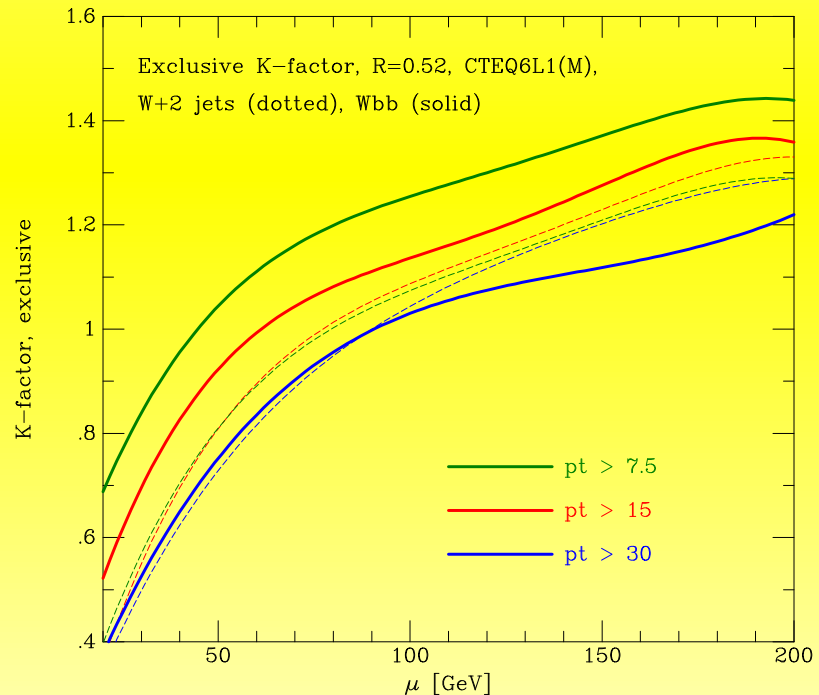
$p_T(\text{jet}) > 7.5 \text{ GeV}$



# Scale dependence of $K$ -factors

- Strong scale dependence.
- The  $Wb\bar{b}$   $K$ -factor varies greatly with the minimum jet  $p_T$ , whereas the  $W + 2$  jets one does not.

dotted:  $W + 2$  jets  
solid:  $W + b\bar{b}$

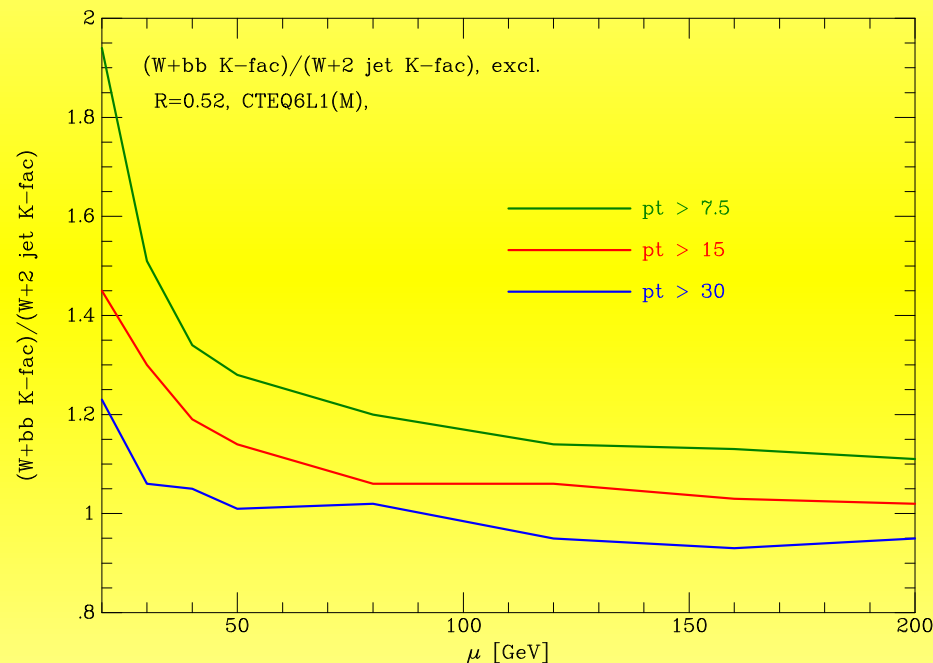


- Scale dependence has a similar shape for both processes.



# *K-factor ratio*

- Important for CDF's "Method 2". Essentially, is a lowest order estimate of  $(Wb\bar{b}/W + 2 \text{ jets})$  reproduced at NLO?

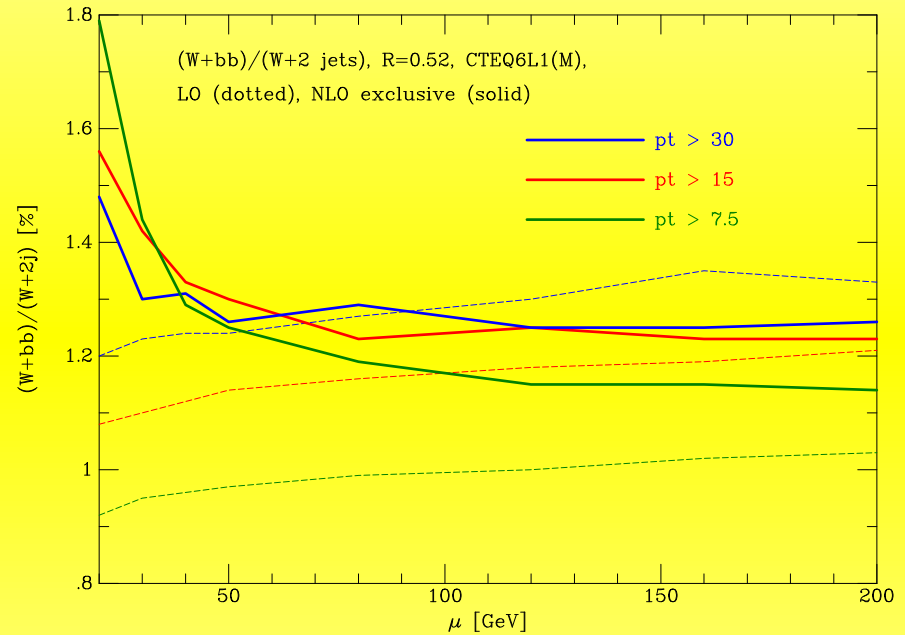


- A qualified "yes" - it is for scale choices around 50 GeV or greater and  $p_T$  cuts of about 15 GeV or greater.
- As the jet  $p_T$  cut is lowered, the ratio gets worse (increases).

# Heavy flavour fraction

- At NLO, this fraction is stable across a wide range of scales.

dotted: LO  
solid: NLO exclusive



- For a  $p_T$  cut of 15 GeV and  $\mu \sim M_W$ , we have:

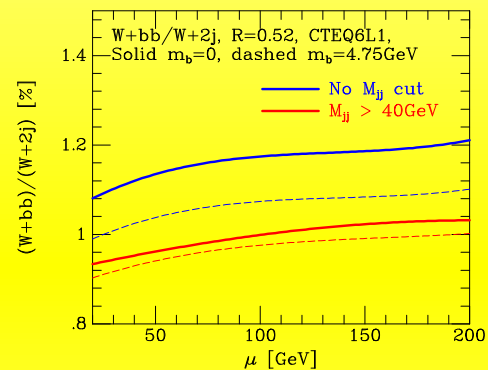
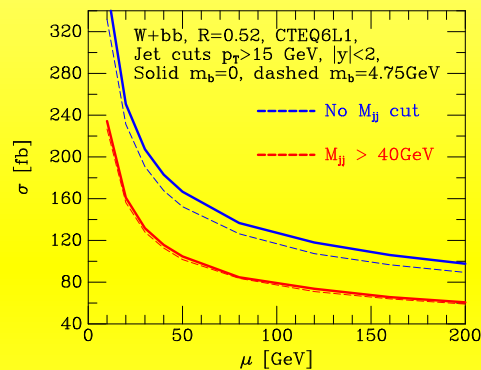
$$\left[ \frac{\sigma(Wb\bar{b})}{\sigma(W + 2 \text{ jets})} \right]_{LO} = 1.16\%,$$

$$\left[ \frac{\sigma(Wb\bar{b})}{\sigma(W + 2 \text{ jets})} \right]_{NLO} = 1.23\%$$

# Effect of NLO corrections on $Wb\bar{b}$ and $W + 2$ jet distributions

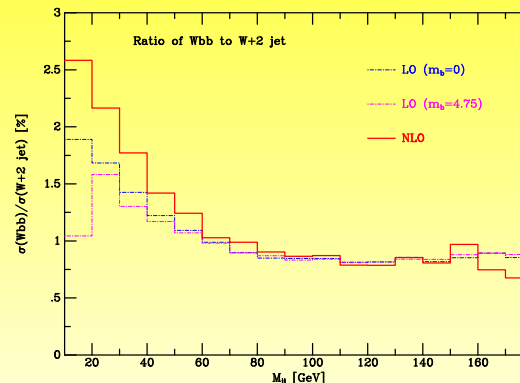
# $b\bar{b}$ mass cut

- Such a cut would be helpful, if it could be experimentally enforced:
  - ★ It improves the massless approximation



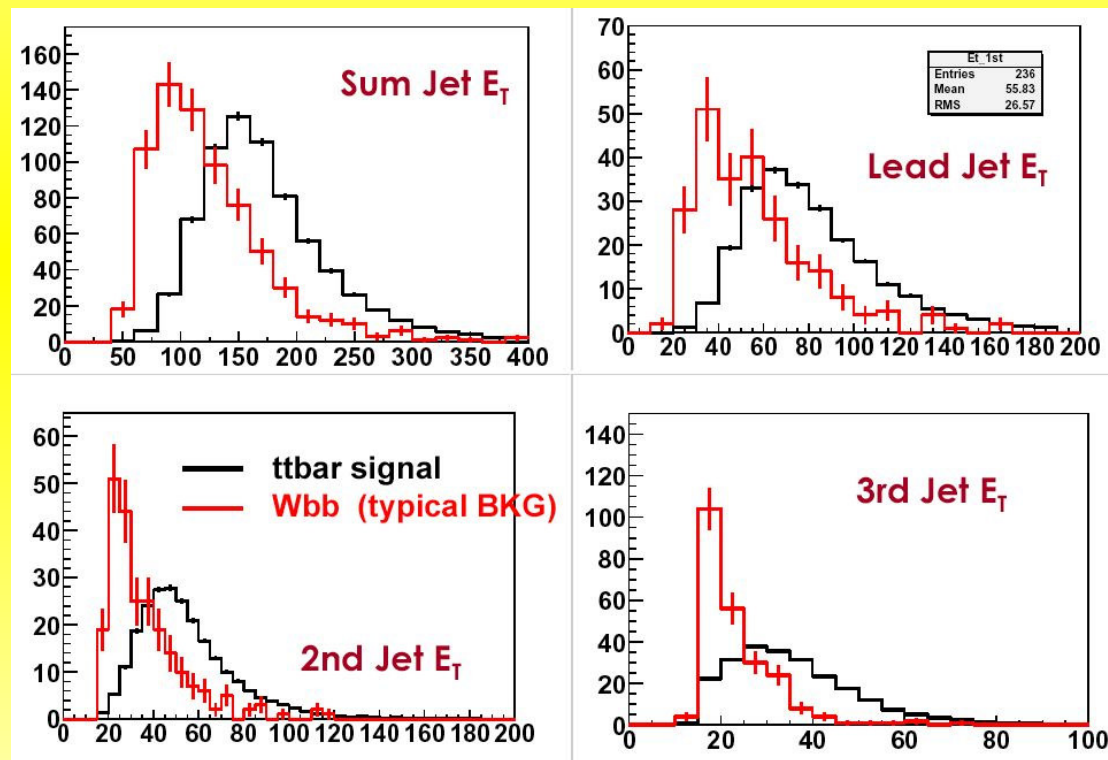
- ★ It reduces this background compared to, for example,  $t\bar{t}$  production, since here the  $b$ 's like to lie at low invariant mass.

LO ( $m_b = 0$ )  
 LO ( $m_b = 4.75$ )  
 NLO ( $m_b = 0$ )



# Shape comparisons

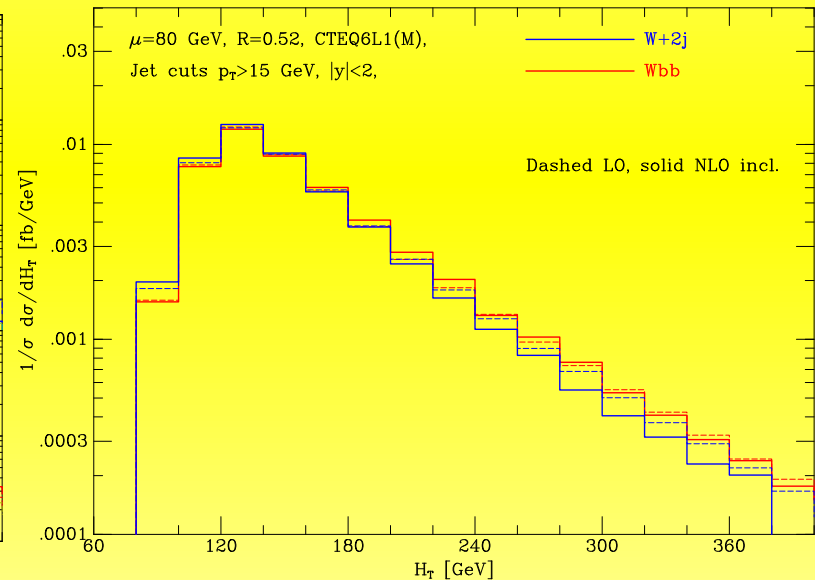
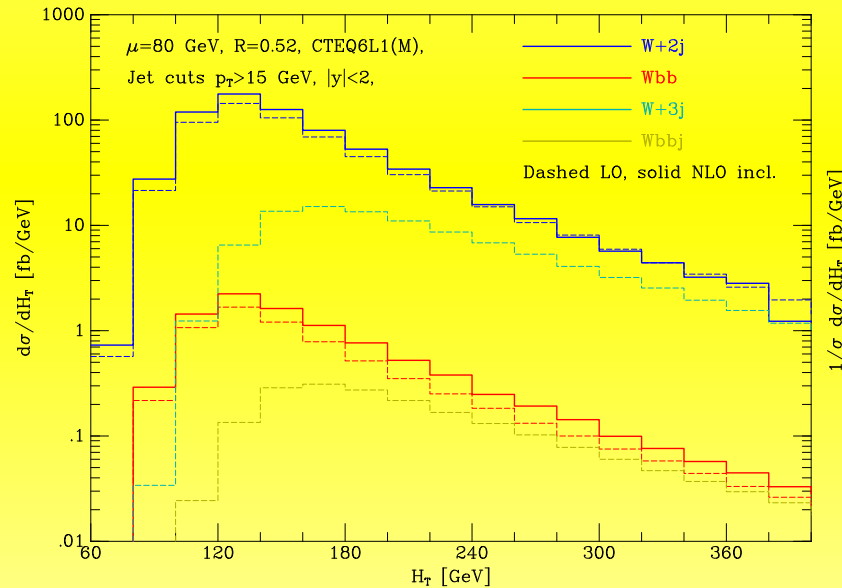
- Shapes of distributions are important in order to either:
  - ★ make kinematic cuts to reduce the  $W + \text{jet}$  backgrounds; or,
  - ★ to fit components of signal and background.



- In particular, are the shapes of relevant distributions similar in the  $b$ -tagged and untagged samples? Is this only true at LO?

# Kinematic distributions

- NLO behaviour may provide clues to processes with more jets ( $\rightarrow$  relevant for  $t\bar{t}$ ), especially for more inclusive variables such as  $\sum E_T(\text{jet})$  and  $H_T = \sum_{\text{event}} E_T$ .



- $Wb\bar{b}$  shape is relatively unchanged at NLO, compared to  $W + 2$  jets.

# NLO predictions

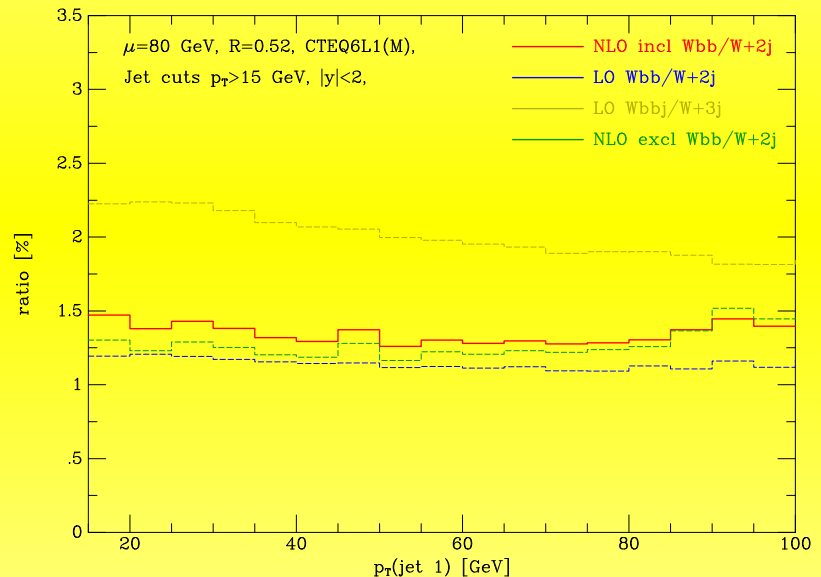
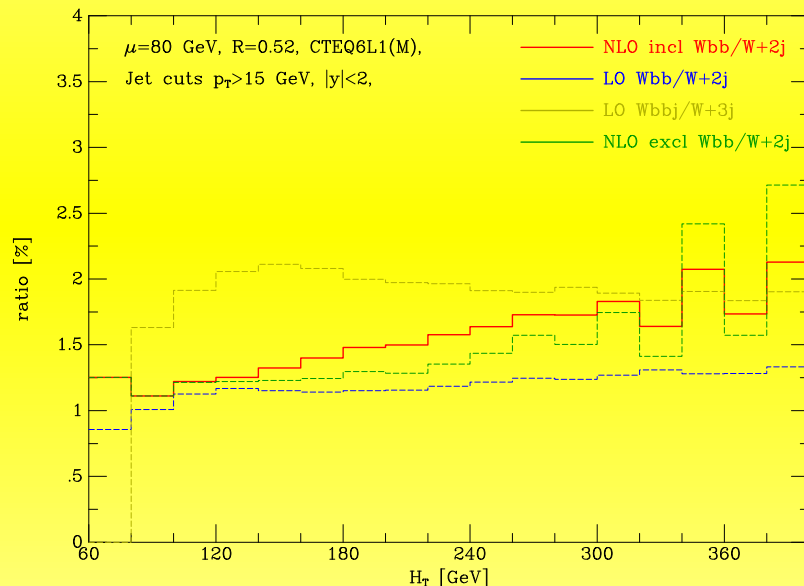
- At NLO, there is a change of shape in the  $H_T$  distribution.

Lowest order

Lowest order+jet

NLO inclusive

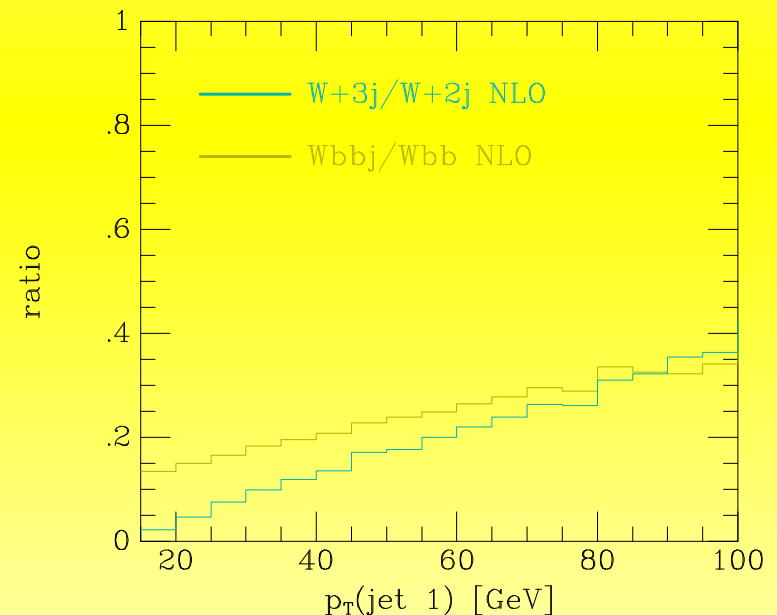
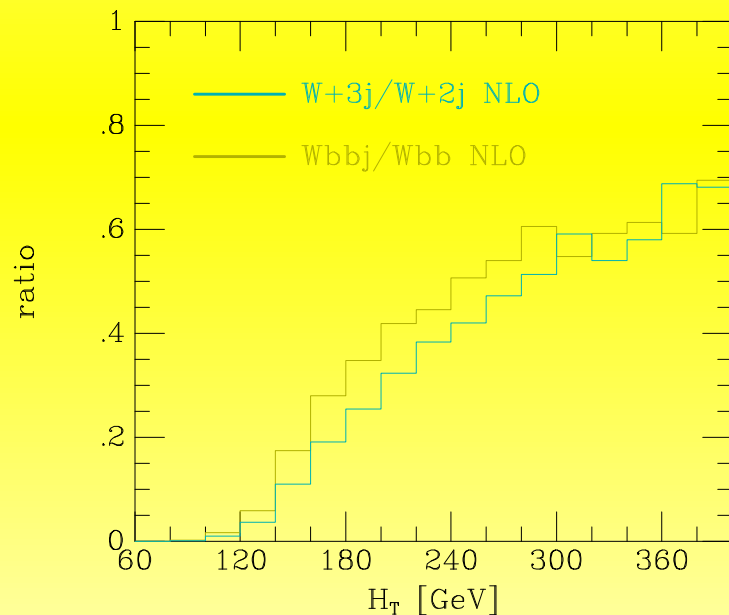
NLO exclusive



- This change is not entirely due to the extra  $W + 3$  jet events allowed in the inclusive sample.
- The  $p_T$  distribution of the hardest jet shows no change in shape.

# Extra jet contribution

- In the NLO inclusive result, the contribution to the  $H_T$  distribution from  $W + 3$  jet events is negligible at small  $H_T$  and dominant at large  $H_T$ .
- Similar ratio for  $Wb\bar{b}j$  to  $Wb\bar{b}$ .

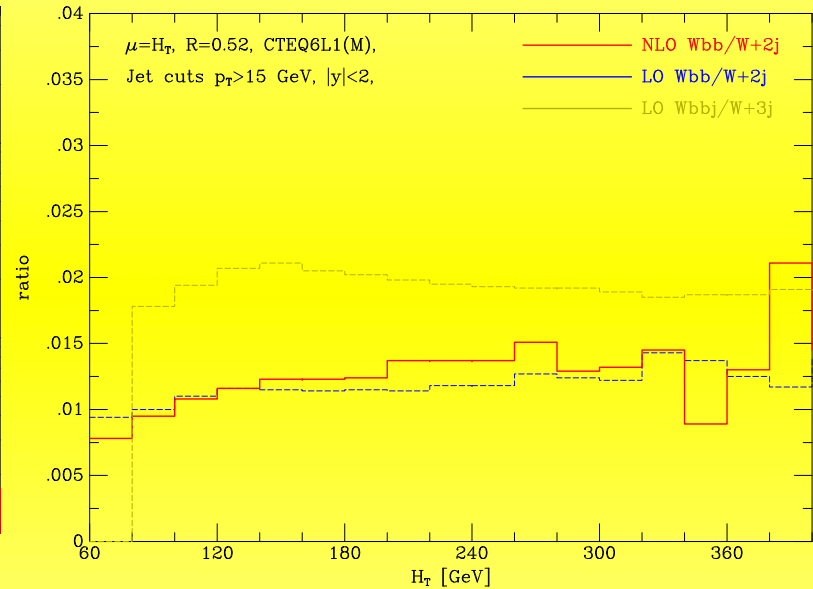
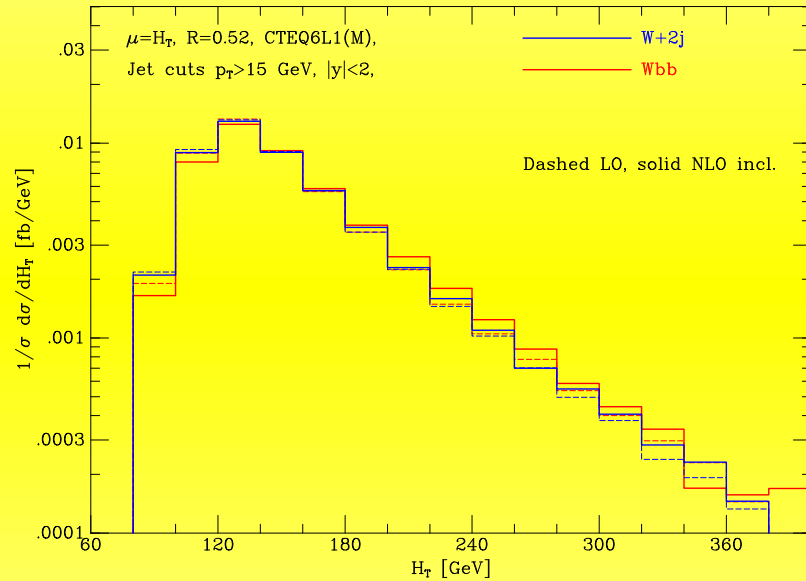


- Extra jet contribution to the jet  $p_T$  distribution is never dominant over this range.



# Dynamical scale

- This behaviour is somewhat altered if a dynamical scale  $\mu = H_T$  is used both at LO and NLO



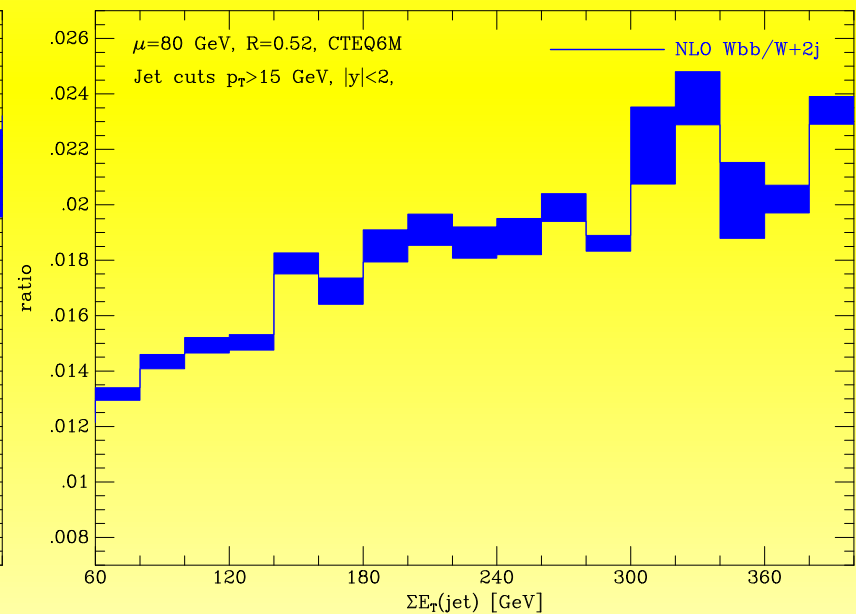
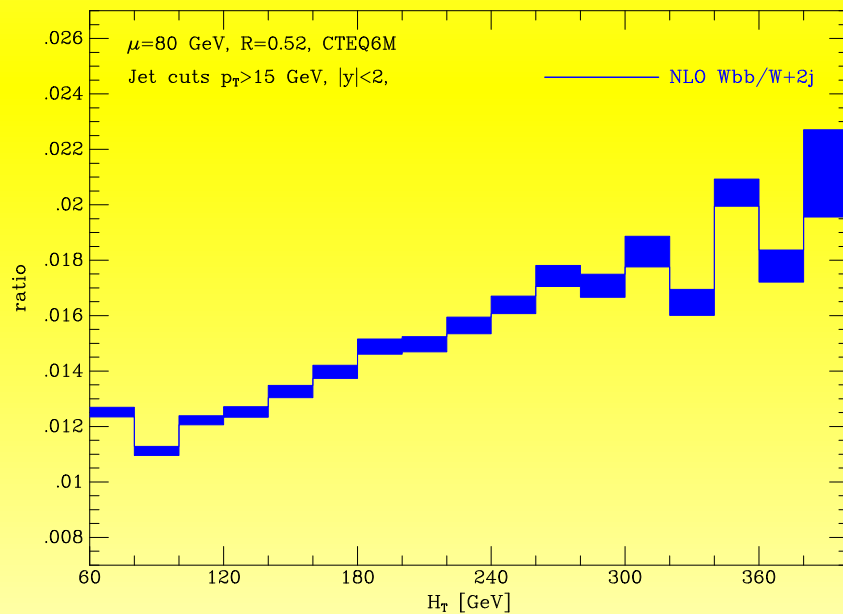
- The shapes of both the  $Wb\bar{b}$  and  $W + 2$  jet distributions change in the same way, leading to a result that is much more encouraging for a LO analysis.

# PDF uncertainties

- Total cross-section uncertainty:

$$Wb\bar{b} \rightarrow 2.5\%, \quad W + 2j \rightarrow 1.5\%.$$

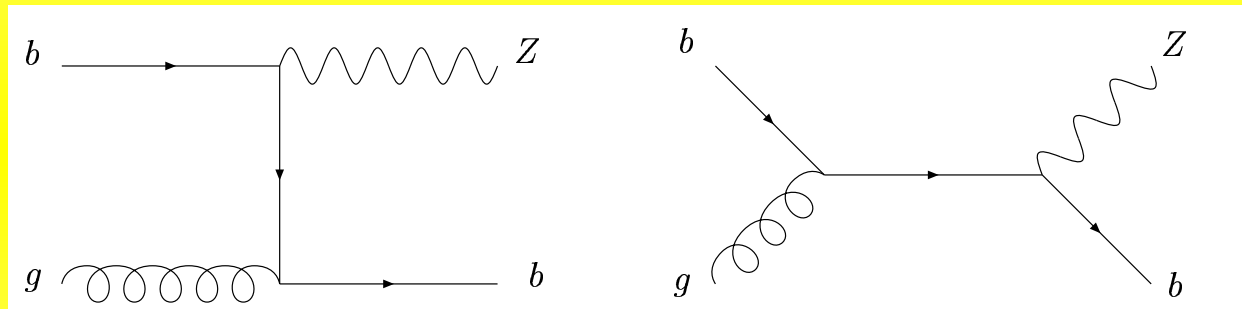
- Uncertainty in the  $(Wb\bar{b}/W + 2 \text{ jet})$  ratio:



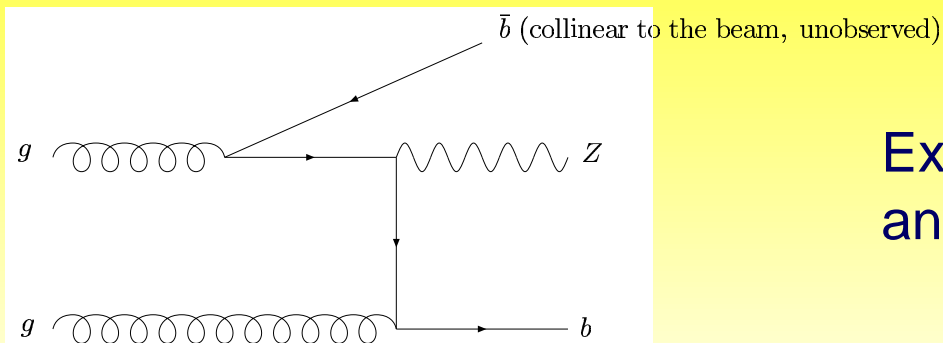
# Single-tagged heavy flavour

# Heavy flavour fraction revisited

- Often the presence of two  $b$ -quarks in the final state is actually only inferred from a single  $b$ -tag
- In this case, there is another way of computing the theoretical cross-section. For instance, in the case of  $Z + \text{heavy flavour}$ :



- Requires knowledge of  $b$ -quark pdf's, but compare to:



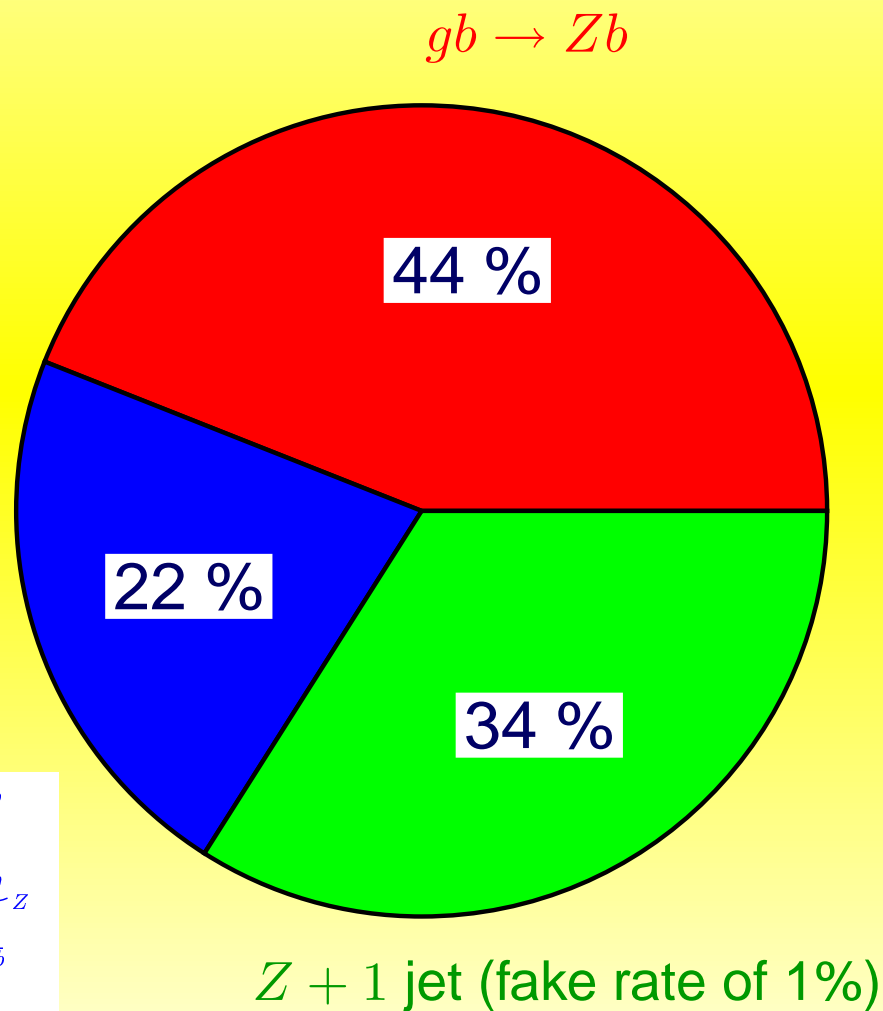
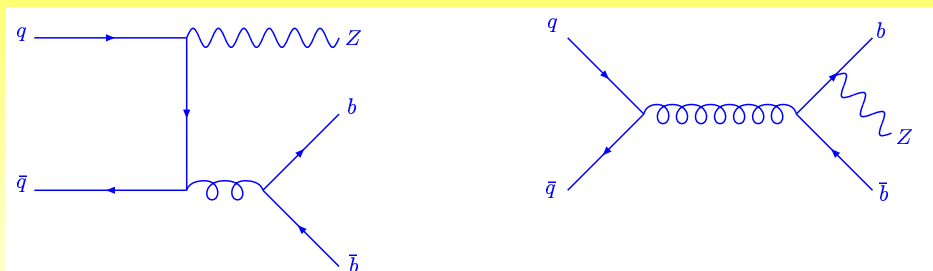
Expansion in  $\alpha_s \ln(M_Z/m_b)$   
and NLO calculation difficult

# $Z + b$ at NLO - Run II

JC, K. Ellis, F. Maltoni and S. Willenbrock, hep-ph/0312024

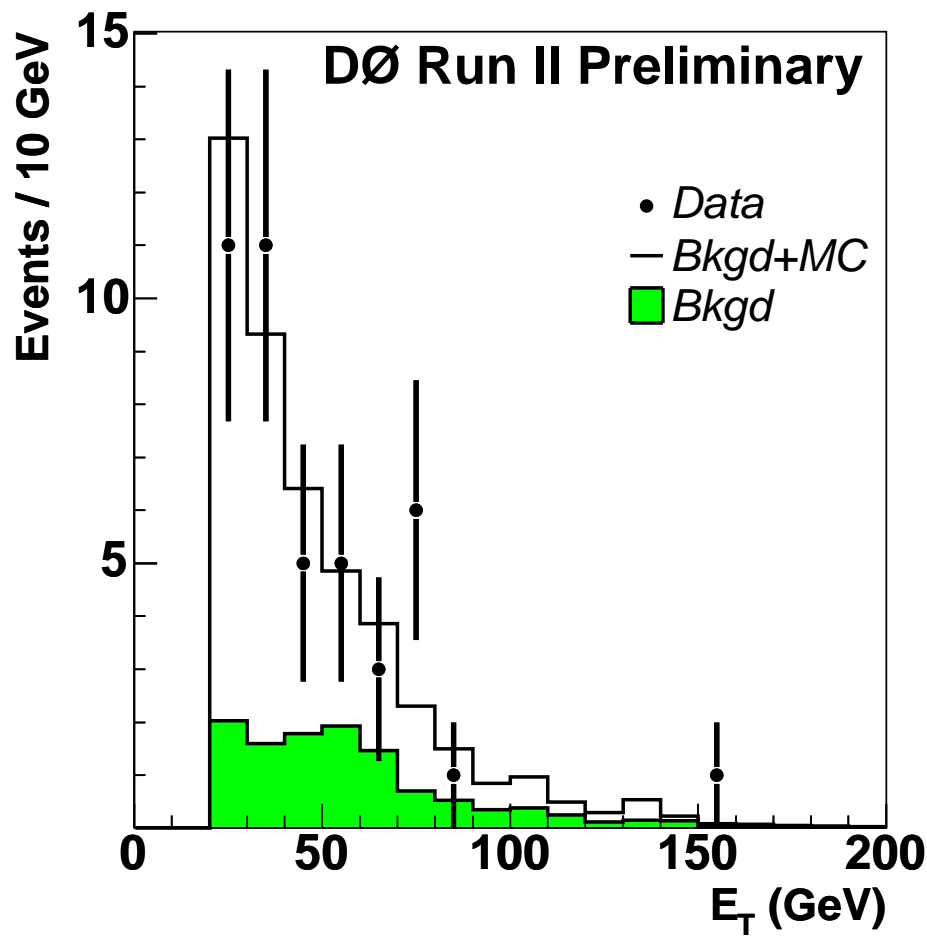
- $p_T^{\text{jet}} > 15 \text{ GeV}, |\eta^{\text{jet}}| < 2$
- $\sigma(Z + \text{one } b \text{ tag}) = 20 \text{ pb}$
- Fakes from  $Z + \text{jet}$  events are significant
- Prediction for ratio of  $Z + b$  to **untagged**  $Z + \text{jet}$  is  $0.02 \pm 0.004$

$$q\bar{q} \rightarrow Z(b\bar{b})$$



# Experimental result

■ Based on  $189 \text{ pb}^{-1}$  of data from Run II



Preliminary ratio of cross-sections:

$$\frac{\sigma(Z+b)}{\sigma(Z+j)} = 0.024 \pm 0.07$$

compatible with the NLO prediction

$Z + b$  process in the next version of MCFM will allow a much better comparison with the analysis

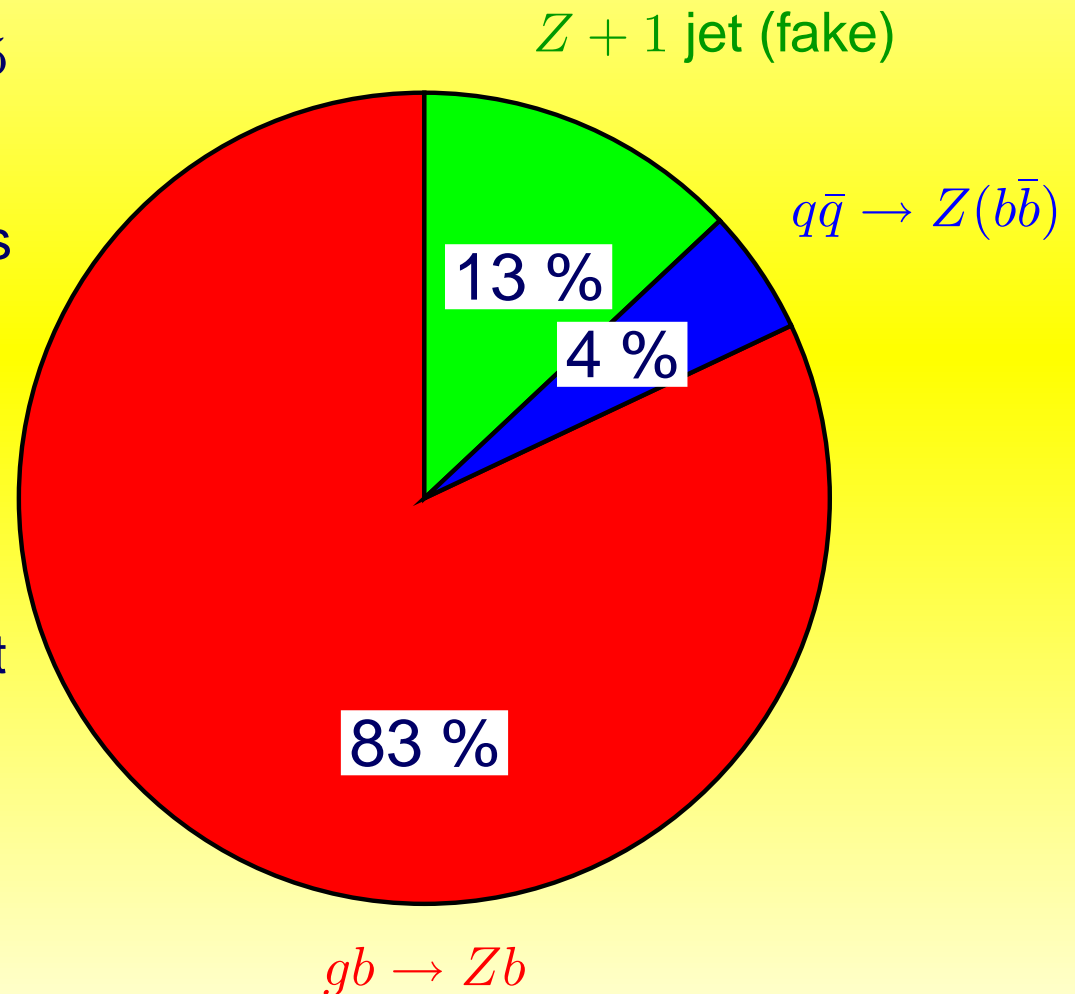
# *LHC expectations*

■  $p_T^{\text{jet}} > 15 \text{ GeV}, |\eta^{\text{jet}}| < 2.5$

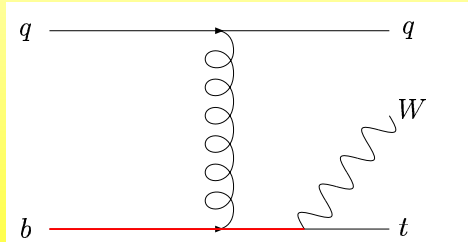
■  $\sigma(Z + \text{one } b \text{ tag}) = 1 \text{ nb}$

■ Fakes from  $Z + \text{jet}$  events are much less significant and  $q\bar{q}$  contribution is tiny

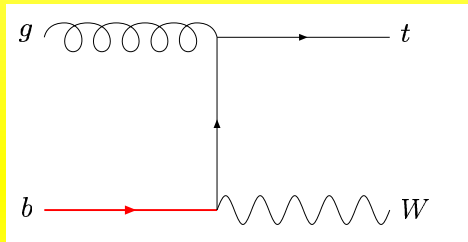
■ This should allow a fairly clean measurement of heavy quark PDF's (currently, only derived perturbatively)



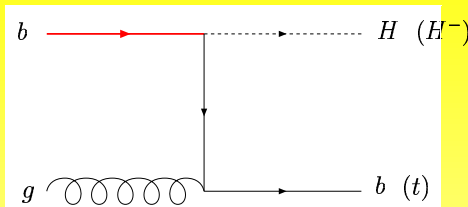
# *b*-PDF uses



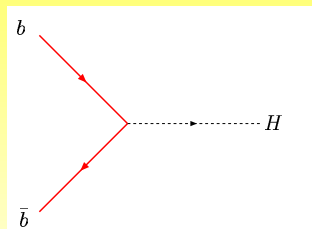
single-top  $q\bar{b} \rightarrow qtW$



single-top  $gb \rightarrow tW$



(charged) Higgs+ $b$



inclusive Higgs



# Computational directions

## ■ $W + 3,4$ jet cross-sections at NLO

- ★ New technology needed: probably not ready for Run II

Nagy and Soper, hep-ph/0308127

Giele and Glover, hep-ph/0402152

## ■ Inclusion of $b$ mass effects in $Wb\bar{b}$ and $Zb\bar{b}$

- ★ Technology available: some efforts are underway ... c.f.  $Hb\bar{b}$

W. Beenakker et al., hep-ph/0211352

S. Dawson et al., hep-ph/0311216

## ■ Merging of existing NLO calculations with a parton shower

- ★ Possible: MC@NLO has yet to be applied to  $W/Z +$  jets

## ■ Further study of recent ideas regarding parton showers (most promising in the short term)

- ★ Matrix elements corrections - CKKW, ...
- ★ How much of the effects of NLO are taken into account by combining matrix element calculations with parton showers?

F. Krauss et al.

# Outlook

- The  $W + \text{jets}$  channel (including heavy quarks) is very important for many studies in Run II.
- Unfortunately, for events with many jets we are limited to LO predictions for rates and distributions.
- However, there should be lots to learn from the NLO corrections that we know about. The highest multiplicity that is currently available is production of  $Wb\bar{b}$  and  $W + 2 \text{ jets}$ .
- Implications for Run II analyses.
  - ★ Results suggest that some relevant observables do not suffer from large NLO effects and we can proceed with more confidence in analyses based on LO tools.
  - ★ However, beware of variables that change shape at NLO ( $H_T$ ).
  - ★ These statements are heavily dependent on scale choices.
- Further comparisons with parton shower approaches and data is the way forward.